

# Spectroscopic studies of $\text{Cu}^{2+}$ ion doped in $\text{CaO-SrO-Na}_2\text{O-B}_2\text{O}_3$ glasses

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Glasses of the  $(20-x)\text{CaO-xSrO-(20-y)Na}_2\text{O-60B}_2\text{O}_3-y$  (CSNB) system with  $(5 \leq x \leq 15)$  mol% and  $y=0.1$ mol% of  $\text{CuO}$  were characterized by X-ray diffraction (XRD), EPR (Electron Paramagnetic Resonance), Optical absorption Spectra and FT-IR Studies. The EPR results indicate that the transition metal ions are present in octahedral site with tetragonal distortion. Spin-Hamiltonian parameters were evaluated. The Optical band energy ( $E_{\text{opt}}$ ) and Urbach energy ( $\Delta E$ ) were calculated from their ultra violet edges. By correlating EPR and Optical data the molecular orbital coefficients have been evaluated. The IR studies shows that the glassy system contains  $\text{BO}_3$  and  $\text{BO}_4$  units in the disordered manner.

(Received March 16, 2012; accepted September 20, 2012)

**Keywords:** EPR, Optical absorption spectra, Borate glasses,  $\text{Cu}^{2+}$  ions, FT-IR

## 1. Introduction

Glasses play an important role in solid state electronic and ionic devices. Borate glasses are of great interest due to their structures and properties which are different from other glasses. Borate glasses are mainly used in electro-optic switches, electro-optic modulators, lasing materials and non-linear optical converters. Many researchers have studied  $\text{B}_2\text{O}_3$  intensively in recent years [1,2], since they are relatively easy to obtain and moreover present interesting structural particularities, due to the existence of the boron anomaly. The boron atom can be placed in the network in tri- or tetra co-ordination depending on the concentration of the modifier oxide [3 – 6].  $\text{B}_2\text{O}_3$  is one of the most important glass forming oxides and has been incorporated into various kinds of glass systems in order to obtain different physical and chemical properties [7]. Transition metal ions are being mostly used in the recent years to find the glass structure, since their outer d-electron orbital function has a broad radial distribution and due to their high sensitive response to the changes in the surrounding actions. Among various transition metal ions, the copper ion, a Para-magnetic metal ion, when dissolved in glass compositions in very small quantities makes the glasses colored and has a strong influence over the optical properties of the glasses [8]. The studies of transition metal ions in glasses by EPR and optical absorption techniques give information on the structure of the glass. In different glasses copper can exist as a monovalent  $\text{Cu}^+$  ion or as divalent  $\text{Cu}^{2+}$  ion to be sensitive to the glass environment. In the present study preparation and characterization of  $(20-x)\text{CaO-xSrO-(20-y)Na}_2\text{O-60B}_2\text{O}_3-y$  (CSNB) glasses were done by using EPR, and optical absorption and IR studies.

## 2. Experimental

The glass samples were prepared by the melt quenching method with the composition given in Table 1. The materials used in the present study are Analar grade  $\text{B}_2\text{O}_3$  (99.9 %),  $\text{CaO}$  (99.9 %),  $\text{SrCO}_3$  (99.9 %) and  $\text{Na}_2\text{CO}_3$  (99.9 %). For transition metal doping 0.1 mol % of  $\text{CuO}$  (99.9 %) is added to the starting materials. These mixtures are sintered at  $500^\circ\text{C}$  and melt in an electric furnace in porcelain crucible at  $900^\circ\text{C}$  -  $1000^\circ\text{C}$  for nearly 1 hr. The melt is then quenched at room temperature in air to form a glass. The glasses so formed are annealed at  $300^\circ\text{C}$  for about 1hr. X-ray powder diffraction patterns of glass samples were recorded using Copper target on Philips PW (1710) diffractometer at room temperature. EPR recordings were made at room temperature on JEOL-JM FE3 with 100 KHz field modulation 100 EPR spectrometer Optical absorption spectra are recorded at room temperature on JASCO V670 Spectrophotometer in 200-900 nm. Infrared transmittance spectra of the powdered glass samples were recorded using JASCO FT-IR 5300 Spectrometer in the wave number range  $400\text{-}4000\text{ cm}^{-1}$  at room temperature.

Table 1. Composition (mol%) of the glasses studied in the present work.

Glass Code	CaO	SrO	$\text{Na}_2\text{O}$	$\text{B}_2\text{O}_3$	CuO
CSNB <sub>0</sub>	15	5	20	60	--
CSNB <sub>1</sub>	5	15	19.9	60	0.1
CSNB <sub>2</sub>	10	10	19.9	60	0.1
CSNB <sub>3</sub>	15	5	19.9	60	0.1

### 3. Results and discussion

#### 3.1 XRD

The X-ray powder diffraction data for two glass systems were obtained from Central Instrumentation Lab, HCU, Hyderabad, India and is shown in Fig. 1. The Philips PW 1710 X-ray powder Diffractometer was used. The results were recorded using a PM 8208A chart recorder and an A100 (Digital) printer with VT125 terminal simultaneously. X-ray diffraction technique was used to check for possible crystallinity of the sample after quenching and annealing. The samples were found to be completely amorphous in nature.

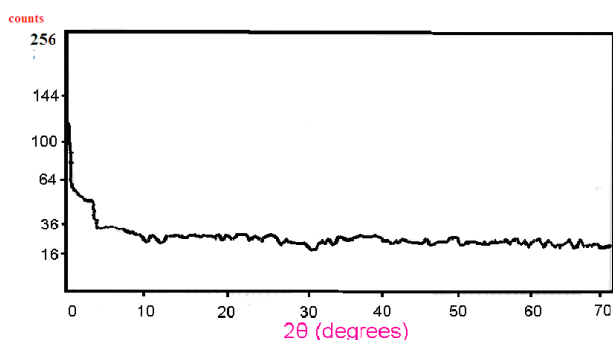


Fig. 1 XRD Spectrum of CSNB glass system.

#### 3.2 EPR study

No EPR signal is observed in undoped glasses confirming that the starting materials used in the present work were free from transition metal impurities. When  $\text{Cu}^{2+}$  ions were added to the CSNB glass systems the EPR spectra exhibited resonance signals similar to those reported for  $\text{Cu}^{2+}$  ions in  $\text{Na}_2\text{O}-\text{B}_2\text{O}_3$ ,  $\text{Na}_2\text{O}-\text{K}_2\text{O}-\text{B}_2\text{O}_3$  glass systems [9, 10]. The  $\text{Cu}^{2+}$  ion with  $S = 1/2$  has a nuclear spin  $I = 3/2$  for both  $^{69}\text{Cu}$  (the natural abundance 69%) and  $^{65}\text{Cu}$  (natural abundance 31%) and therefore  $(2I+1)$  i.e., four parallel and four perpendicular hyperfine components would be expected. In the observed spectra, three parallel components were observed in the lower field region. The perpendicular components in the higher field region were well resolved. The EPR spectra of the three glass systems doped with  $\text{Cu}^{2+}$  ions are shown in Fig. 2. The spectrum closely resembles that of the  $\text{Cu}^{2+}$  ion in most oxide glasses [11,12] from the spectral analysis the spin-Hamiltonian parameters are calculated using the equation [13] and are presented in Table 2. From Table 2, it is clear that  $g_{\parallel} > g_{\perp} > g_e$  i.e.,  $\text{Cu}^{2+}$  is in an octahedral coordination with tetrahedral distortion and the ground

state of  $\text{Cu}^{2+}$  is  $d_{x^2-y^2}$ . Also  $g_{\parallel}$  and  $g_{\perp}$ ,  $A_{\parallel}$  and  $A_{\perp}$  ( $B_1=40$ ,  $B_2=46$ ,  $B_3=53$ ) values increases with the composition by varying CaO and SrO content in CSNB glasses systems and  $g_{\parallel}$  reaches a maximum for  $x = 15\%$  SrO, and  $x = 5\%$  CaO and then reaches a minimum for  $x = 10\%$  CaO and  $x = 10\%$  SrO in CSNB glass system. Thus from Fig.2, Table 2, it is clear that the EPR spectra of  $(20-x)\text{CaO}-x\text{SrO}-(20-y)\text{Na}_2\text{O}-60\text{B}_2\text{O}_3-y\text{CuO}$  glasses systems were strongly concentration dependent. The EPR spectra were analyzed using the spin-Hamiltonian

$$H = g_{\parallel} \beta H_z S_z + g_{\perp} \beta (H_x S_x + H_y S_y) + A_{\parallel} S_z I_z + A_{\perp} (S_x I_x + S_y I_y) \quad (1)$$

Here Z is the symmetry axis,  $\beta$ , is the Bohr magneton, S and I are the electron and nuclear spin operators,  $H_x$ ,  $H_y$  and  $H_z$  are the static magnetic field components,  $g_{\parallel}$  and  $g_{\perp}$  are the parallel and perpendicular components of the hyperfine tensor A. The values of  $A_{\parallel}$  are calculated using the following equation due to Kivelson [14] from Table 2. It is clear that  $g_{\parallel}$  and  $g_{\perp}$  values increases slightly and reach a minimum in CSNB glasses at about 15% CaO.  $g_{\parallel}$  and  $g_{\perp}$  values increase abruptly and reaches a maximum at about 10% CaO and 10% SrO. It is also clear from Table 2, that  $A_{\parallel}$  value decreases with decrease of CaO and there is considerable increase in  $A_{\perp}$  at about 5% CaO and 15% SrO. This indicates that the distortion around  $\text{Cu}^{2+}$  ion is changing with the decreasing content of CaO and increasing content of SrO, and also indicates a continuous structural change in the glass systems.

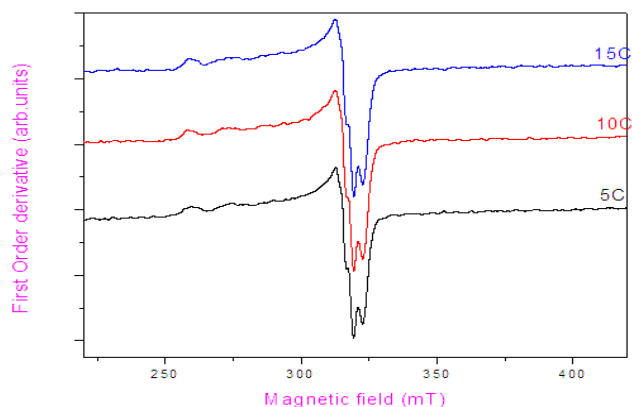


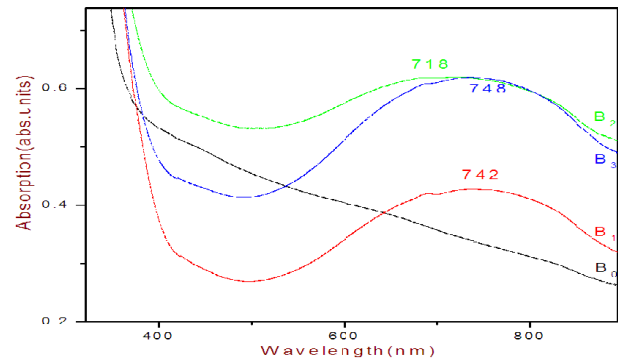
Fig. 2. EPR Spectrum of  $\text{Cu}^{2+}$  doped CSNB glass systems.

Table 2. Comparison of spin-Hamiltonian parameters of Cu<sup>2+</sup> ions in different systems.

System	g <sub>  </sub>	g <sub>⊥</sub>	A <sub>  </sub> x 10 <sup>-4</sup> (Cm <sup>-1</sup> )	Reference
Na <sub>2</sub> O- B <sub>2</sub> O <sub>3</sub>	2.327	2.065	150	9
Na <sub>2</sub> O- K <sub>2</sub> O-B <sub>2</sub> O <sub>3</sub>	2.338	2.046	135	10
CaO-SrO-Na <sub>2</sub> O- B <sub>2</sub> O <sub>3</sub> (1)	2.329	2.066	117	Present work
CaO-SrO-Na <sub>2</sub> O- B <sub>2</sub> O <sub>3</sub> (2)	2.335	2.055	123	Present work
CaO-SrO-Na <sub>2</sub> O- B <sub>2</sub> O <sub>3</sub> (3)	2.351	2.056	130	Present work

### 3.3 Optical absorption spectrum

Optical absorption spectra for all glass composition exhibit a single broad peak (Fig. 3) in the range 715-750 nm, which can be attributed to the presence of Cu<sup>2+</sup> ion in the glass [15]. This absorption can be assigned to the <sup>2</sup>E<sub>g</sub>(D)→<sup>2</sup>T<sub>2g</sub>(D) transition of Cu<sup>2+</sup> ion, by varying CaO and SrO content the optical absorption peak range shifted towards higher wavelength region. The absorption peak position of CSNB glasses at different concentration of CaO and SrO are given in Table 3, similar results were reported in literature [16-19]. Experimental data shows that the cupric ion generally exists in solutions solids and glasses in octahedral symmetry with a strong tetragonal distortion [16-20]. According to the present EPR studies Cu<sup>2+</sup> ions in CSNB glasses present in octahedral symmetry with elongated tetragonal distortion.

Fig. 3. Optical bands of Cu<sup>2+</sup> doped CSNB glass systems.Table 3. Absorption peak of Cu<sup>2+</sup> bonding parameters and % of bonding symmetry of Cu<sup>2+</sup> doped in CaO-SrO-Na<sub>2</sub>O- B<sub>2</sub>O<sub>3</sub> glasses.

Sample	Cu <sup>2+</sup> peak (nm)	Γ <sub>σ</sub> (%)	Γ <sub>π</sub> (%)	α <sup>2</sup>	β <sub>1</sub> <sup>2</sup>	β	K
CSNB <sub>1</sub>	742	61.22	16	0.719042857	0.923148195	0.960806013	0.3703
CSNB <sub>2</sub>	718	57.09	12	0.737995237	0.948542021	0.973931219	0.3458
CSNB <sub>3</sub>	748	57.52	10	0.736014285	0.952889383	0.976160531	0.3337

### 3.4 Optical basicity

The Optical basicity of an oxide glass will reflect the ability of glass to donate negative charge to the probe ion [21]. Duffy and Ingram [22] proposed that the optical basicity can be predicted from the composition of the glass and basicity moderating parameters of various cations present. The theoretical values of the optical basicity (Δ<sub>th</sub>)

of the glasses can be calculated using the following formula,

$$\Delta_{th} = \sum_{i=1}^n \frac{Z_i r_i}{|Z_o| \gamma_i} \quad (2)$$

where ‘n’ is the number of cations present ‘Z’<sub>i</sub> is the oxidation number of the i<sup>th</sup> cation, ‘r<sub>i</sub>’ is the ratio of number of i<sup>th</sup> cation to the number of oxides present and ‘γ<sub>i</sub>’ is the basicity moderating parameters, of the i<sup>th</sup> cation. The basicity moderating parameters ‘γ<sub>i</sub>’ can be calculated from the following equation.

$$\gamma_1 = 1.36 (x_i - 0.26) \tag{3}$$

where ‘x<sub>i</sub>’ is Pauling electro negativity of the cation. The theoretical values of the optical basicity (Δ<sub>th</sub>) of all the glass samples were calculated and are presented in Table 4. It is observed that with the increase SrO and decrease of CaO content the value of optical basicity increases.

Table 4. Summary of Direct, Indirect, Urbach energies and Optical Basicity.

Sample	Direct	Indirect	Urbach	Optical basicity
CSNB <sub>0</sub>	3.968	3.882	0.257	0.5374
CSNB <sub>1</sub>	3.763	3.724	0.269	0.5369
CSNB <sub>2</sub>	3.749	3.699	0.270	0.5388
CSNB <sub>3</sub>	3.719	3.690	0.271	0.5405

### 3.5 Optical energy and Urbach energy

The values of the optical energy gap and Urbach energy calculated in the present work for different glasses are presented in Table 4 and is shown in Fig. 4,5,6. The values obtained in the present work are of the same order for those of copper tellurium oxide [23] and borate glasses [24]. From the values of the widths of the tails of the localized states (ΔE) within the optical band gap for the present glasses, it can be observed that the direct and indirect values increases and Urbach energy values decreases. The decrease in Urbach energy with Cu<sup>2+</sup> concentration can be considered as due to decreased defects [25].

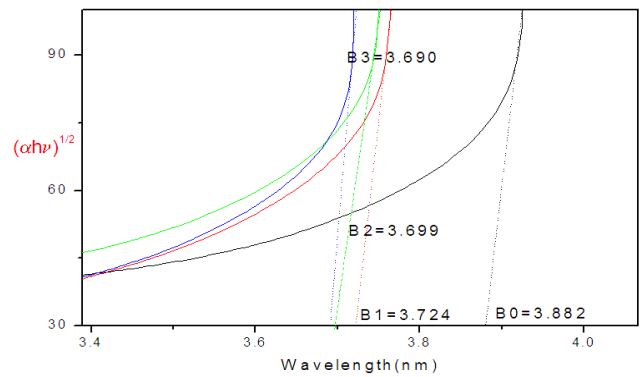


Fig. 4. Indirect bands of Cu<sup>2+</sup> doped CSNB glass systems.

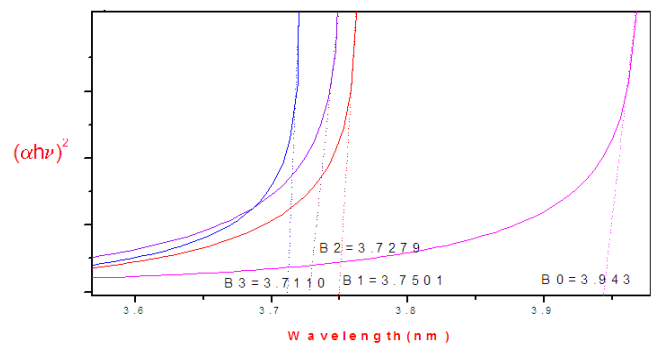


Fig. 5. Direct bands Cu<sup>2+</sup> doped CSNB glass systems.

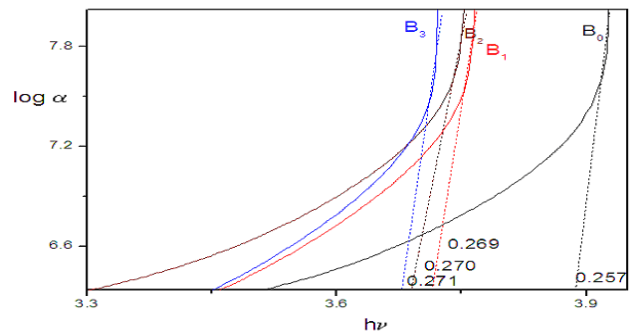


Fig. 6. Urbach energy bands of Cu<sup>2+</sup> doped CSNB glass systems.

### 3.6 Molecular Orbital coefficients

The EPR and optical absorption data can be correlated to evaluate the bonding coefficients as follows [19]

$$g_{||} = 2.0023 \left[ 1 - \frac{4 \lambda \alpha^2 \beta_1^2}{E_1} \right] \tag{4}$$

$$g_{\perp} = 2.0023 \left[ 1 - \frac{\lambda \alpha^2 \beta_1^2}{E_2} \right] \tag{5}$$

where  $\lambda$ , spin orbit coupling parameter is equal to  $-828 \text{ cm}^{-1}$  for CuO and  $\beta^2 \approx 1$  for octahedral environment.  $\Delta E_{xy}$  ( $E_1$ ) and  $\Delta E_{xz,yz}$  ( $E_2$ ) are heights of the  $d_{xy}$  and  $d_{xz,yz}$  and molecular orbital levels above the ground state  $d_{x^2-y^2}$  and these values are estimated from optical absorption spectra [26]. In optical absorption spectra the position of observed absorption maximum of Cu<sup>2+</sup> ions indicates the value of  $\Delta E_{xy}$  from Table 3. It is observed that the bonding parameters are changing with the percent of CaO and SrO. The bonding co-efficient  $\alpha^2, \beta_1^2$  and  $\beta$  characterize respectively. The in plane  $\sigma$  bonding in-plane  $\Pi$  bonding and out of plane  $\Pi$  bonding of the copper (II) complex their values lie between 0.5 and 1.0 the limits of pure covalent and pure ionic bonding [27]. The value of the calculated  $\alpha^2, \beta_1^2$  indicate that the in plane  $\sigma$  bonding is covalent. The normalized covalency of Cu-O in-plane bonding of  $\sigma$  and  $\Pi$  symmetry are expressed by [28] in terms of bonding coefficients  $\alpha^2, \beta_1^2$  as

$$\Gamma_{\sigma} = \frac{200(1-S)(1-\alpha^2)}{1-2S} \% \quad (6)$$

$$\Gamma_{\Pi} = 200(1-\beta_1^2) \% \quad (7)$$

where  $s$  is the overlap integral ( $S_{oxy} = 0.076$ ). The normalized covalency values of the Cu (II) – O of in-plane bonding of  $\Pi$  symmetry ( $\Gamma_{\Pi}$ ). The calculated are given in Table 3. The bonding coefficient  $\alpha^2$  (in-plane  $\sigma$  bonding) can be calculated from the EPR data using the following expression given by Kuska et al [29].

$$\alpha^2 = \frac{7}{4} \left[ \frac{A_{\parallel}}{P} - \frac{A}{P} - \frac{2}{3} g_{\parallel} - \frac{5}{21} g_{\perp} + \frac{6}{7} \right] \quad (8)$$

where  $P = 0.036 \text{ cm}^{-1}$  and  $A = (1/3A_{\parallel} + 2/3A_{\perp})$  the  $\alpha^2$  value calculated from the above equation was used in equation 10 and 11 to evaluate  $\beta_1^2$  from the spin-Hamiltonian parameters the dipole term ( $P$ ) and the Fermi-contact term ( $K$ ) are calculated using the expressions [30, 31].

$$P = 2\gamma_{Cu} \beta_0 \beta_N (r^{-3}) = 0.036 \text{ cm}^{-1} \quad (9)$$

$$K = (A_0/P) + \Delta g_0 \quad (10)$$

Here  $\gamma_{Cu}$  is the magnetic moment of copper,  $\beta_0$  is the Bohr magneton,  $\beta_N$  is the nuclear magneton and  $r$  is the distance from the central nucleus to the electron.  $A_0 = (A_{\parallel} + 2A_{\perp})/3$ , where  $A_{\parallel}$  and  $A_{\perp}$  are the hyperfine coupling constants in the parallel and perpendicular

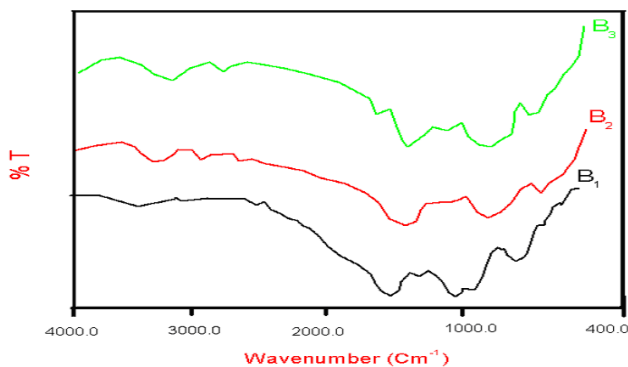
directions to the field.  $g_{\parallel}$  and  $g_{\perp}$  are the  $g$ -values parallel and perpendicular to the field and  $\Delta g_0 = g_0 - g_e$ , where  $g_0 = (g_{\parallel} + 2g_{\perp})/3$  and  $g_e$  is the free ion ( $g_e = 2.0023$ ). The Fermi contact term  $K$  is a measure of the polarization produced by the uneven distribution of  $d$ -electron density on the inner core  $s$ -electron. The evaluated values of  $K$  (Table 3) are tune with the general order of [32].

### 3.7 FT- IR studies

Usually the vibrational modes of the borate network are found [33] to be active in three infrared spectral regions, which are similar to those reported by several workers. First group bands was reported to occur at  $1200-1600 \text{ cm}^{-1}$  and is due to the asymmetric stretching relaxation of the B-O band of trigonal  $\text{BO}_3$  units, whereas second group lies between  $800$  and  $1200 \text{ cm}^{-1}$  and is due to the B-O band stretching of the tetrahedral  $\text{BO}_4$  units. Finally the third group is observed around  $700 \text{ cm}^{-1}$  and is attributed to bending of B-O-B linkages in the borate networks. The principal frequencies observed in the IR spectrum of the glasses studied and their assignments are summarized in Table 5 and is shown in Fig. 7. The spectra corresponding to glasses of systems containing the absorption bands of the vitreous  $\text{B}_2\text{O}_3$  There is the band at  $\sim 700-713 \text{ cm}^{-1}$  assigned to the B-O-B deformation vibrational mode [34, 35]. The band at  $\sim 951 \text{ cm}^{-1}$  characteristic of glasses with high  $\text{B}_2\text{O}_3$  content is due to boroxol rings [36]. Its intensity diminished, when  $X = 10\%$  CaO and  $X = 10\%$  SrO content. The band at  $1006 - 1016 \text{ cm}^{-1}$  is assigned to the B-O stretching mode including tetra-coordinated boron atoms ( $\text{BO}_4$ ). The intensity of this band increase when  $X = 15\%$  SrO concentration increases. Around  $1207-1249 \text{ cm}^{-1}$  a band assigned to the B-O stretching including three coordinated boron atoms ( $\text{BO}_3$ ), [34-38]. By comparing the amplitude of bands at  $\sim 1016 \text{ cm}^{-1}$  and  $\sim 1249 \text{ cm}^{-1}$ . The dominant presence of the  $\text{BO}_4$  units as compared to the  $\text{BO}_3$  units is obvious. Around  $\sim 1360 \text{ cm}^{-1}$  a shoulder belonging to the  $1249 \text{ cm}^{-1}$  band of the boroxol rings was identified. The appearance of the H-O-H bending mode centered at  $1726 \text{ cm}^{-1}$  proves the hygroscopic character of the glasses investigated. The intensity of this band increases when  $X = 5\%$  CaO and  $X = 15\%$  SrO content of the sample. The absorption band at around  $2847-2920 \text{ cm}^{-1}$  can be attributed to the hydrogen bonds. The same peaks have been recorded [33]. In the  $3410-3431 \text{ cm}^{-1}$  spectral a broad band assigned to the H-O stretching mode (hydrogen bonded) was identified. The band around  $1410-1419 \text{ cm}^{-1}$  containing diborate groups Very weak absorption band is observed at  $472 \text{ cm}^{-1}$  when  $X = 5\%$  CaO tentatively assigned to  $\text{CaO}_4$  tetrahedra on the basis of a vibrational study of various crystalline calcium compounds  $\text{CaO}_4$  tetrahedra give bands in the  $400 - 500 \text{ cm}^{-1}$  region. The observations made in the present investigation agree with the literature values [8, 39-41]. The structural changes may be explained by the partial replacement of oxygen atoms by varying CaO and SrO content.

Table 5. FT-IR studies.

Wave number (Cm <sup>-1</sup> )			Assignment
CSNB <sub>1</sub>	CSNB <sub>2</sub>	CSNB <sub>3</sub>	
472	--	--	Ca-O Stretching mode
609	--	--	Sr-O Stretching mode
713	700	717	B-O-B deformation Vibrational mode
951	--	951	Boroxol ring vibration
1014	1016	1006	B-O stretching mode of BO <sub>4</sub> units
1149	--	--	B-O stretching mode of BO <sub>3</sub> units
--	--	1207	B-O stretching mode of BO <sub>3</sub> units
--	--	1249	B-O stretching mode of BO <sub>3</sub> units
1419	1410	1415	B-O-B deformation mode
--	--	1726	H-O-H bending mode
--	2847	2851	H-O-H bending mode
--	2918	2920	H-O-H bending mode

Fig. 7. FT-IR Spectra of Cu<sup>2+</sup> doped CSNB glass systems.

#### 4. Conclusions

1. From the EPR and Optical spectra of Cu<sup>2+</sup> ions in CSNB glasses, it is confirmed that the copper ions occupied octahedral sites with tetragonal distortion.
2. The structural changes take place with decreasing CaO and increasing SrO composition.
3. The optical absorption spectra of these samples give a single broad band due to Cu<sup>2+</sup> ions in distorted octahedral sites.

4. FT-IR studies show that the glassy system contains BO<sub>3</sub> and BO<sub>4</sub> units in a disordered manner.
5. The optical band gaps, direct, indirect, and Urbach energies are found to vary with CaO and SrO content, which is assigned to the increase in defect concentration with CuO content.
6. By correlating EPR and optical data, the molecular orbital coefficients  $\alpha^2, \beta_1^2$  are evaluated. The values suggest that the in-plane  $\sigma$  bonding is moderately covalent in nature.

#### Acknowledgement

One of the authors (J. Lakshmi Kumari, letter No. F. ETFA PANO. 84 / FIP-XI Plan) is thankful to UGC, New Delhi for providing financial assistance under FDP Programme to carry out these investigations.

Dr. Sandhya Cole would like to thank the University Grants Commission, New Delhi for sanctioning UGC-MRP F. No. 9-498/2010(SR).

#### References

- [1] Eskisehir, G. U. Journal of Science **22**(3), 129 (2009).
- [2] G. Lakshminarayana, S. Buddhudu, Spectrochim, Acta, part A, **63**(2), 295 (2006).

- [3] I. Tawansi, A. Gohar, D. Holland, N. A. E-I. Shishatawi, *J. Phys and Apply. phys.* **21**, 607 (1988).
- [4] H. Ushida, Y. Iwadate, T. Hattori, S. Nishiyama, K. Fukushima, Y. Ikeda, M. Yamaguchi, M. Misawa, T. Fukunaga, T. Nakazawa, S. Jitsukawa, *J. Alloys Compounds* **377**, 167 (2004).
- [5] T. Takaishi, J. Jin, M. Takaishi, T. Uchino, T. Yoko. *Proc. Int. Congr. Glass. 2. Extended Abstracts. Edinburgh. Scotland, 1-6* (2001).
- [6] I. Yasui, H. Hasegawa, Y. Saito, Y. Akasaka, *J. Non-Cryst. Solids* **123**, 71 (1990).
- [7] T. Yana, N. Kunimine, S. Shibata, M. Yamane, *J. Non-Cryst. Solids* **321**, 147 (2003).
- [8] G. Upender, V. Kamalaker, C. P. Vardhani, V. Chandra Mouli, *Indian Journal of pure and Applied Physics* **47**, 551 (2009).
- [9] R. V. S. S. N. Ravi Kumar, V. Raja gopal Reddy, A. V. Chandra Sekhar, B. J. Reddy, Y. P. Reddy, P. S. Rao, *Journal of Alloys and Compounds* **337**, 272 (2002).
- [10] R. P. Sreekanth Chakradhar, K. P. Ramesh, J. L. Rao, J. Rama Krishna, *J. Phys. Condens, Matter* **15**, 1469 (2003).
- [11] S. Suresh, J. Chinna Babu, V. Chandra Mouli, *Phys. Chem (Italic) Glasses* **46**, 27 (2005).
- [12] J. D. Lee *Concise inorganic chemistry* (Black well Science, Oxford) (1996).
- [13] J. M. Dance, J. Drmaudery, H. Baudry, M. Moneraya, *Solid State Comm* **39**, 199 (1981).
- [14] D. Kivelson, Lee-SK, *J. Chem. phys* **47**, 11 (1986).
- [15] S. Sakka, K. Kamira, K. Makita, Yamanoto, *J. Non-Cryst. Solids* **63**, 223 (1984).
- [16] T. Bates, in "Modern Aspects of the Vitreous State", edited by J. D. Mackenzie (Butter Worths, London) **2**, 195 (1962).
- [17] R. Harani, C. A. Hogarth, K. A. K. Lott. *J. Mater. Sci* **19**, 1420 (1984).
- [18] I. Ardelean, O. Cozar, S. Filip, V. Pop, I. Cenan *Solid State Commun* **100**, 609 (1996).
- [19] R. R. Kumar, A. K. Bhatnagar, B. C. V. Reddy, *Solid State Commun* **114**, 493 (2000).
- [20] A. Klonkowski, *Phy. Chem. Glasses* **24**, 116 (1983).
- [21] A. J. Easteal, A. T. Marcom, *J. Non-cryst. Solids* **34**, 29 (1979).
- [22] J. A. Duffy, M. D. Ingram, *J. Inorg. Nul. Chem.* **37**, 1203 (1975).
- [23] M. A. Hassan, C. A. Hogarth, *J. Mater Sci.* **23**, 2500 (1988).
- [24] A. I. Sabry, M. M. E-Samanoudy, *J. Mater Sci* **30**, 3930 (1995).
- [25] M. Vithal, P. Nachimuthu, T. Banu, R. Jagannathan, *J. Apply. Phys.* **81**, 7922 (1997).
- [26] Krosh-Moe, *J. phys. Chem. Glasses* **6**, 46 (1965).
- [27] A. Klonkowski, *Phys. Chem. Glasses* **24**, 116 (1983).
- [28] H. Kawazoe, H. Hosono, T. Kanazawa, *J. Non Cryst Solids* **33**, 103 (1979).
- [29] H. A. Kuska, M. T. Rogers, R. E. Durringer, *J. Phys. Chem.* **71**, 109 (1967).
- [30] D. Kivelson, R. Neiman, *J. Chem. phys* **35**, 149 (1961).
- [31] K. E. Folk, E. Ivanova, B. Roos, T. Vanogard, *Inorg. chem.* **9**, 556 (1970).
- [32] J. H. Vanveleck, *Phys. Rev.* **41**, 208 (1932).
- [33] S. G. Motke, S. P. Yawale, S. S. Yawale, *Bull Mater. Sci.* **25**, 75 (2002).
- [34] F. F. Bentley, L. D. Smithson, A. L. Rozee, *Infrared Spectra and Characteristic Frequencies 700- 300 cm<sup>-1</sup>*.
- [35] B. Sharma, D. Dubc, A. Mansingh, *J. Non-Cryst. Solids* **69**, 39 (1984).
- [36] C. S. Sundandana, R. Singh, *Phys. Status Solidi (A)* **K91** (1983).
- [37] E. Culea, I. Bratu, A. L. Nicula, *ibid*, **83**, K 15 (1984).
- [38] C. N. R. Rao "Chemical Applications of Infrared spectroscopy" (Academic press. Newyork) (1963).
- [39] I. Ardelean, M. Peteanu, R. Ciceo-Lucacel, *Journal of Material Science, Materials in Electronics* **11**, 11 (2000).
- [40] S. Suresh, M. Prasad, G. Upender, V. Kamalaker, V. Chandra Mouli, *Indian Journal of pure and Applied Physics* **47**, 163 (2009).
- [41] N. Srinivasa Rao, M. Purnima, Shashidhar, Bale, K. Siva Kumar, Syed Rahman, *Bull. Mater Sci* **29**, 365 (2006).

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